

Dynamical spike solutions in a nonlocal model of pattern formation

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Joint work with

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Nonlocal Models and Peridynamics

Technische Universität, Berlin

November 4, 2012

Reaction-diffusion equations

The point of departure:

a general system of reaction-diffusion (**Ordinary-PDE**) equations:

$$\begin{aligned}u_t &= f(u, v), & \text{for } x \in \bar{\Omega}, t > 0 \\v_t &= D\Delta v + g(u, v) & \text{for } x \in \Omega, t > 0\end{aligned}$$

in a bounded domain $\Omega \subset \mathbb{R}^n$.

The Neumann boundary condition:

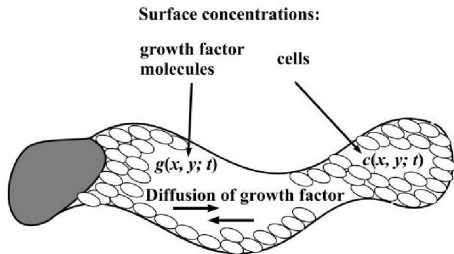
$$\partial_n v = 0 \quad \text{for } x \in \partial\Omega, t > 0$$

Initial data:

$$u(x, 0) = u_0(x), \quad v(x, 0) = v_0(x).$$

- $D > 0$ – a constant diffusion coefficient.
- arbitrary C^1 -nonlinearities $f = f(u, v)$ and $g = g(u, v)$.

Biological system



- Cell proliferation is influenced by growth factor
- Growth factor is externally supplied or produced by the cells
- Growth factor diffuses along the structure formed by the cells and binds to cell membrane receptors
- **Hypothesis:** The diffusion of this growth factor may significantly influence the dynamics of the whole cell population

Derivation of "SHADOW SYSTEM"

Lemma

Let $T > 0$. Assume that $v^D = v^D(x, t)$ are solutions for each $D > 0$ of

$$v_t = D\Delta v + F$$

$$\partial_n v = 0$$

$$v(x, 0) = v_0(x).$$

in a bounded domain $\Omega \subset \mathbb{R}^n$ and for $t \in [0, T]$.

Some nice assumptions of u_0 and $F = F(x, t)$.

Then, for every $t > 0$, the limit

$$\lim_{D \rightarrow \infty} v^D(\cdot, t) = \xi(t)$$

exists where

$$\frac{d}{dt} \xi(t) = \int_{\Omega} F(x, t) dx, \quad \xi(0) = \int_{\Omega} v_0(x) dx.$$

Idea of the proof

First, we consider the solution $v^D = v^D(x, t)$ of

$$v_t = D\Delta v$$

$$\partial_n v = 0$$

$$v(x, 0) = v_0(x).$$

Then $v^D(\cdot, t)$ converges towards a constant function

$$\lim_{D \rightarrow \infty} v^D(\cdot, t) \rightarrow \int_{\Omega} v_0(x) dx \equiv \xi(0) \quad \text{each } t > 0.$$

Hint: rescale the time and use the identity $v^D(x, t) = v^1(x, Dt)$.

Idea of the proof

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For the inhomogeneous problem, we use the Duhamel principle

$$v(x, t) = e^{tD\Delta} v_0(x) + \int_0^t e^{(t-s)D\Delta} F(x, s) ds.$$

In the limit, we obtain the function

$$\xi(t) = \int_{\Omega} v_0(x) dx + \int_0^t \int_{\Omega} F(x, s) dx ds.$$

Shadow problem

The solutions $(u^D, v^D) = (u^D(x, t), v^D(x, t))$ of ordinary-PDE problem

$$\begin{aligned}u_t &= f(u, v), \\v_t &= D\Delta v + g(u, v)\end{aligned}$$

converge as $D \rightarrow \infty$ towards a solution

$$(u, \xi) = (u(x, t), \xi(t))$$

of the following **shadow system**

$$\begin{aligned}u_t &= f(u, \xi), & \text{for } x \in \bar{\Omega}, t > 0 \\ \xi_t &= \int_{\Omega} g(u(x, t), \xi(t)) dx & \text{for } t > 0\end{aligned}$$

in a bounded domain $\Omega \subset \mathbb{R}^n$ supplemented with initial data

$$u(x, 0) = u_0(x), \quad \xi(0) = \xi_0.$$

Local existence of solutions

Shadow problem

$$\begin{aligned}u_t &= f(u, \xi), & \text{for } x \in \overline{\Omega}, t > 0 \\ \xi_t &= \int_{\Omega} g(u(x, t), \xi(t)) dx & \text{for } t > 0\end{aligned}$$

with initial data: $u(x, 0) = u_0(x), \quad \xi(0) = \xi_0.$

Local existence of solutions by standard methods via the integral equations

$$\begin{aligned}u(x, t) &= u_0(x) + \int_0^t f(u(x, s), v(x, s)) dx \\ \xi(t) &= \xi_0 + \int_0^t \int_{\Omega} g(u(x, s), \xi(s)) dx ds.\end{aligned}$$

Stationary solutions

Shadow system

$$\begin{aligned}u_t &= f(u, \xi), \\ \xi_t &= \int_{\Omega} g(u(x, t), \xi(t)) dx\end{aligned}$$

We calculate a stationary solution $(U(x), \xi_0)$ from equations

$$f(U(x), \xi_0) = 0, \quad \int_{\Omega} g(U(x), \xi_0) dx = 0.$$

If the equation $f(U(x), \xi_0) = 0$ has a unique solution, we obtain **only constant stationary solutions**.

Remark

One can also consider non-constant stationary solutions. They have to be constant on subsets of Ω .

Instability of steady states

Theorem

Assume that the constant vector $(\bar{u}, \bar{\xi})$ is a solution of the initial-boundary value problem for the system

$$\begin{aligned}u_t &= f(u, \xi), \\ \xi_t &= \int_{\Omega} g(u(x, t), \xi(t)) \, dx,\end{aligned}$$

which means that $f(\bar{u}, \bar{\xi}) = 0$ and $g(\bar{u}, \bar{\xi}) = 0$. If

$$f_u(\bar{u}, \bar{\xi}) > 0,$$

then $(\bar{u}, \bar{\xi})$ is *unstable solution* of this initial-boundary value problem.

- The inequality $f_u(\bar{u}, \bar{\xi}) > 0$ can be interpreted as **the autocatalysis of u at the steady state $(\bar{u}, \bar{\xi})$** .
- It appears in a natural way in models describing **diffusion-driven instability**.

Idea of the proof

Step 1. Linear instability.

The linearized system at the steady state $(\bar{u}, \bar{\xi})$:

$$U_t = f_u(\bar{u}, \bar{\xi})U + f_\xi(\bar{u}, \bar{\xi})\Phi$$

$$\Phi_t = g_u(\bar{u}, \bar{\xi}) \int_{\Omega} U \, dx + g_\xi(\bar{u}, \bar{\xi})\Phi$$

The unbounded solution is given explicitly

$$U(x, t) = e^{tf_u(\bar{u}, \bar{\xi})} U_0(x), \quad \Phi(t) \equiv 0,$$

where $U_0(x)$ is any function such that $\int_{\Omega} U_0(x) \, dx = 0$.

Step 2. Nonlinear instability.

One should adapt classical methods from reaction-diffusion equations.

Model example

We explain Instability theorem in the case of the problem

$$u_t = -u + u^2 \xi, \quad \text{for } x \in \bar{\Omega}, t > 0$$

$$\xi_t = -\xi - k\xi \int_{\Omega} u^2(x, t) dx + B \quad \text{for } t > 0$$

$$u(x, 0) = u_0(x), \quad \xi(0) = \xi_0$$

where $k, B \in \mathbb{R}$ are fixed positive parameters. Here,

$$g(u, \xi) = -\xi - k\xi u^2 + B.$$

For simplicity of notation, we assume that $|\Omega| = 1$.

Space homogeneous solutions

Theorem

All solutions $(u(t), \xi(t))$ of the following initial value problem for ordinary differential equations

$$\begin{aligned}\frac{d}{dt}u &= -u + u^2\xi, & \frac{d}{dt}\xi &= -\xi - ku^2\xi + B \\ u(0) &= u_0 \geq 0, & \xi(0) &= \xi_0 \geq 0\end{aligned}$$

are nonnegative, *global-in-time*, and *uniformly bounded for $t > 0$* .

Proof.

We observe that

$$\frac{d}{dt}(ku(t) + \xi(t)) = -(ku(t) + \xi(t)) + B.$$

Hence, as long as $u(t)$ and $\xi(t)$ are nonnegative, they have to be uniformly bounded for $t > 0$.

Space homogeneous solutions

Constant stationary solutions

We solve the system of equations

$$-u + u^2\xi = 0, \quad -\xi - ku^2\xi + B = 0$$

So,

$$\text{either } u = 0, \quad \xi = B \quad \text{or} \quad u = \frac{1}{\xi}, \quad -\xi - \frac{k}{\xi} + B = 0$$

Remark

The trivial solution $(0, B)$ is an asymptotically stable solution of ODE.

Remark

Two nontrivial positive solutions exist if $B > 0$ is large enough.

One of them is stable as a solution of ODE. Another one is unstable.

Shadow problem

Remark

The trivial solution $(0, B)$ is also an asymptotically stable solution of the shadow problem.

Remark

Autocatalysis condition:

For $f_u(u, \xi) = -1 + 2u\xi$ at the nontrivial constant solution $u = 1/\xi$ we have

$$f_u(u, \xi) = 1 > 0.$$

Nontrivial constant solutions are unstable as solutions of the shadow problem.

This is an example of [the Diffusion-Driven Instability](#) or [the Turing instability](#) in the case of the shadow problem.

Question

Nonnegative **space homogeneous** solutions of the shadow problem

$$u_t = -u + u^2\xi, \quad \text{for } x \in \bar{\Omega}, t > 0$$

$$\xi_t = -\xi - k\xi \int_{\Omega} u^2(x, t) dx + B \quad \text{for } t > 0$$

$$u(x, 0) = u_0(x), \quad \xi(0) = \xi_0$$

are bounded and global-in-time.

What is the large time behavior of space inhomogeneous solutions ?

Question

Nonnegative **space homogeneous** solutions of the shadow problem

$$u_t = -u + u^2 \xi, \quad \text{for } x \in \bar{\Omega}, t > 0$$

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$$u(x, 0) = u_0(x), \quad \xi(0) = \xi_0$$

are bounded and global-in-time.

What is the large time behavior of space inhomogeneous solutions ?

Some space inhomogeneous solutions of this shadow problem may blow up in finite time.

Blowup of solutions

Theorem

Let $0 \in \Omega$ and assume that $u_0 \in C(\Omega)$ satisfies:

$$u_0(0) = 1 \quad \text{and} \quad 0 \leq u_0(x) < 1 \quad \text{for} \quad x \neq 0$$

and

$$A_0 \equiv \int_{\Omega} \left(\frac{u_0(x)}{1 - u_0(x)} \right)^2 dx < \infty.$$

Assume also that

$$\min \left\{ \xi_0, \frac{B}{1 + kA_0} \right\} > 1.$$

Then, the corresponding solution of shadow problem for the system

$$u_t = -u + u^2 \xi, \quad \xi_t = -\xi - k\xi \int_{\Omega} u^2(x, t) dx + B$$

blows up in a finite time at $x = 0$.

Idea of the proof.

For fixed $\xi(t)$ and for each $x \in \Omega$, we solve the equation $u_t = -u + u^2\xi$:

$$u(x, t) = \frac{e^{-t}}{\frac{1}{u_0(x)} - \int_0^t \xi(s)e^{-s} ds}.$$

Note that

$$T_{max} = \sup \left\{ t > 0 : \int_0^t \xi(s)e^{-s} ds < 1 \right\}$$

because

$$u_0(0) = 1 \quad \text{and} \quad 0 \leq u_0(x) < 1 \quad \text{for} \quad x \neq 0$$

Hence, we have an estimate up to the blowup point:

$$u(x, t) \leq \frac{e^{-t}}{\frac{1}{u_0(x)} - 1} = \frac{u_0(x)e^{-t}}{1 - u_0(x)} \quad \text{for all} \quad (x, t) \in \Omega \times [0, T_{max}).$$

Idea of the proof.

Next, using the estimate of $u(x, t)$ we deduce from the equation for ξ the following differential inequality

$$\xi_t \geq -(1 + kA_0)\xi + B \quad \text{for all } t \in [0, T_{max})$$

which implies the lower bound

$$\xi(t) \geq \min \left\{ \xi_0, \frac{B}{1 + kA_0} \right\} \quad \text{for all } t \in [0, T_{max}).$$

Choosing parameters correctly, we obtain the blowup.

Activator-inhibitor system

Gierer and Meinhardt system:

$$u_t = \varepsilon \Delta u - u + \frac{u^p}{v^q},$$
$$\tau v_t = D \Delta v - v + \frac{u^r}{v^s},$$

with the Neumann boundary conditions and with positive initial data.

One assumes that $\varepsilon > 0$ is very small and $D > 0$ is large (comparing with the domain size).

In the two extreme cases, we have the shadow problem

$$u_t = -u + \frac{u^p}{\xi^q},$$
$$\tau \xi_t = -\xi + \int_{\Omega} \frac{u^r}{\xi^s} dx.$$

Caution: In the literature, the expression “shadow problem” is used in the case of the Gierer–Meinhardt system with $D = +\infty$ and $\varepsilon > 0$.

Activator-inhibitor system

$$u_t = -u + \frac{u^p}{\xi^q},$$

$$\tau \xi_t = -\xi + \int_{\Omega} \frac{u^r}{\xi^s} dx.$$

We assume that $|\Omega| = 1$.

Results

- The vector $(1, 1)$ is a constant stationary solution.

Activator-inhibitor system

$$u_t = -u + \frac{u^p}{\xi^q},$$
$$\tau \xi_t = -\xi + \int_{\Omega} \frac{u^r}{\xi^s} dx.$$

We assume that $|\Omega| = 1$.

Results

- The vector $(1, 1)$ is a constant stationary solution.
- It is asymptotically stable as a solution of ODE, provided $\tau > 0$ is small.

Activator-inhibitor system

$$u_t = -u + \frac{u^p}{\xi^q},$$
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Results

- The vector $(1, 1)$ is a constant stationary solution.
- It is asymptotically stable as a solution of ODE, provided $\tau > 0$ is small.
- The vector $(1, 1)$ is **unstable** as a solution of the shadow problem.

Activator-inhibitor system

$$u_t = -u + \frac{u^p}{\xi^q},$$
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We assume that $|\Omega| = 1$.

Results

- The vector $(1, 1)$ is a constant stationary solution.
- It is asymptotically stable as a solution of ODE, provided $\tau > 0$ is small.
- The vector $(1, 1)$ is **unstable** as a solution of the shadow problem.
- For some range of parameters p, q, r, s , solutions of the ODE system are bounded and global-in-time.

Activator-inhibitor system

$$u_t = -u + \frac{u^p}{\xi^q},$$
$$\tau \xi_t = -\xi + \int_{\Omega} \frac{u^r}{\xi^s} dx.$$

We assume that $|\Omega| = 1$.

Results

- The vector $(1, 1)$ is a constant stationary solution.
- It is asymptotically stable as a solution of ODE, provided $\tau > 0$ is small.
- The vector $(1, 1)$ is **unstable** as a solution of the shadow problem.
- For some range of parameters p, q, r, s , solutions of the ODE system are bounded and global-in-time.
- For the same range of parameters, some space inhomogeneous solutions **blow up** in finite time.

Model of early carcinogenesis

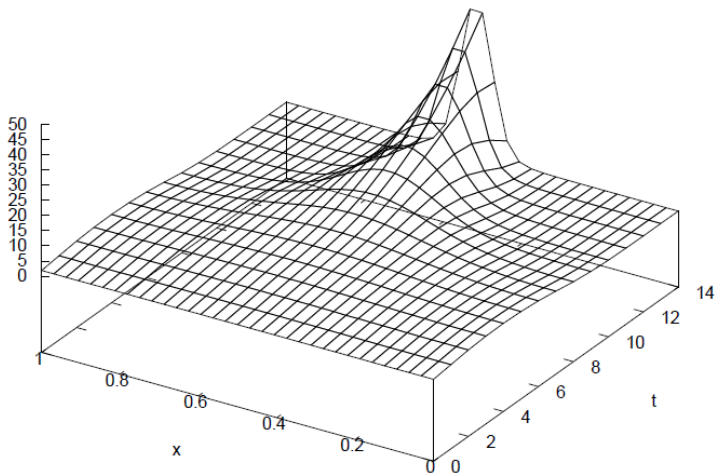
$$u_t = \left(\frac{au\xi}{1+u\xi} - 1 \right) u$$
$$\xi_t = -\xi - k\xi \int_{\Omega} u^2 dx + B$$

in a bounded set $\Omega \subset \mathbb{R}^n$. We complete this system with nonnegative initial conditions

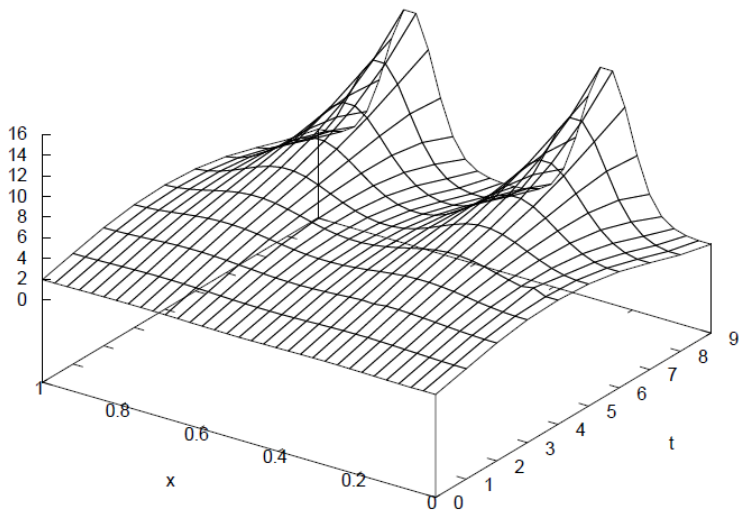
$$u(0, x) = u_0(x), \quad \xi(0) = \xi_0.$$

- Here, nonnegative solutions are **global-in-time**.
- Blowup of solutions appears when $t \rightarrow \infty$.

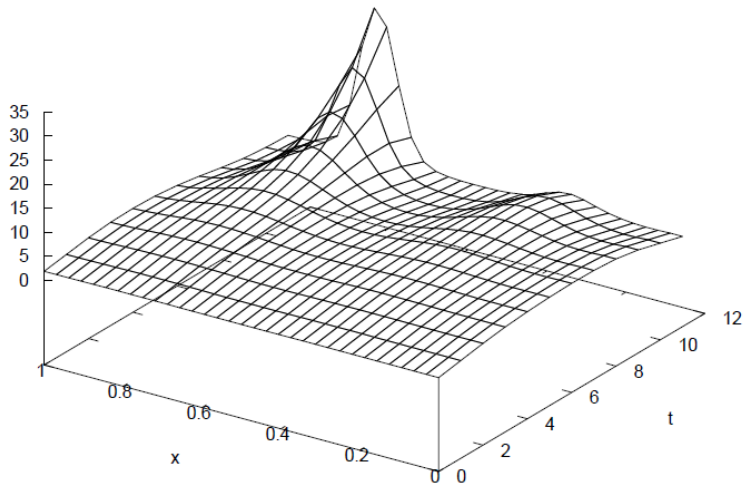
One point blowup



Two points blowup



One point blowup



Initial datum which lead to blowup

Assumption for blowup:

Let $0 \in \Omega$ and assume that $u_0 \in C(\Omega)$ satisfies:

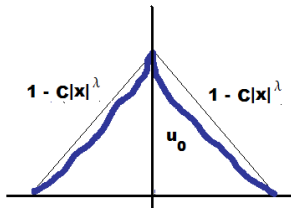
$$u_0(0) = 1 \quad \text{and} \quad 0 \leq u_0(x) < 1 \quad \text{for} \quad x \neq 0$$

and

$$A_0 \equiv \int_{\Omega} \left(\frac{u_0(x)}{1 - u_0(x)} \right)^2 dx < \infty.$$

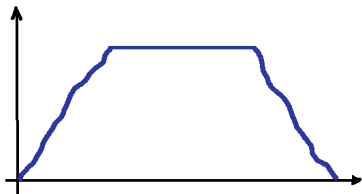
Remark

The constant A_0 is finite if there exist constants $C > 0$ and $\lambda \in (0, n/2)$ such that $u_0(x) \leq 1 - C|x|^\lambda$ for all $x \in \Omega$.



Initial datum with NO blowup

The following initial datum cannot generate blowup.



For the system

$$u_t = -u + u^2 \xi, \quad \text{for } x \in \bar{\Omega}, t > 0$$

$$\xi_t = -\xi - k\xi \int_{\Omega} u^2(x, t) dx + B \quad \text{for } t > 0$$

$$u(x, 0) = u_0(x), \quad \xi(0) = \xi_0$$

the "mass" is bounded:

$$\sup_{t \in [0, T_{\max})} \int_{\Omega} u(x, t) dx < +\infty.$$

THANK YOU