

# Existence, uniqueness and stability of semi-linear rough partial differential equations

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## Rough PDE via Feynman-Kac integrals

Rough PDEs: motivation from stoch. filtering

Rough linear PDE

Rough semilinear PDE

## Rough PDE via Davies' method

Motivation: Burgers equation

Variational approach

Rough differential equations according to Davie

A priori estimates

Stability

# Motivation: Stochastic filtering theory

Estimating partially observed diffusion process

$$(S) \quad dX_t = b(X_t) dt + \sigma(X_t) dB_t \quad \in \quad \mathbb{R}^d$$

$$(O) \quad dW_t = \gamma(X_t) dt + d\tilde{B}_t \quad \in \quad \mathbb{R}^e$$

for (ind.) Brownian motions  $B, \tilde{B}$

cond. expectation of (S) given  $\mathcal{W}_{s:t} = \sigma(W_u : u \in [s, t])$

$$E^{t,x} [g(X_T) | W_{t,T}] = \frac{E_0^{t,x} \left[ g(X_T) \exp \left( \int_t^T \gamma(X_s) dW_s - \frac{1}{2} \int_t^T \|\gamma(X_s)\|^2 ds \right) \right]}{E_0^{t,x} \left[ \exp \left( \int_t^T \gamma(X_s) dW_s - \frac{1}{2} \int_t^T \|\gamma(X_s)\|^2 ds \right) \right]}$$

where  $E_0^{t,x}$  denotes expectation w.r.t.  $B$  only (!)

## Rem

- ▶ crucial pb: robustness w.r.t. data  $W_{0:t}$ , (in part.  $\int_t^T \gamma(X_s) dW_s$ )
- ▶ **our Ansatz** consider the rough path topology

## Rough path integrals - main idea

- ▶ recall Young's ineq  $|\int_s^t Y_r - Y_s dX_r| \leq C_{\beta,\alpha} |t - s|^{\alpha+\beta} \|Y\|_\beta \|X\|_\alpha$ ,  $\alpha + \beta > 1$
- ▶ but for  $Y_t = F(X_t)$ ,  $F \in C^2$

$$\begin{aligned} \int_s^t F(X_r) dX_r &= F(X_s) \underbrace{(X_t - X_s)}_{=: X_{st}} + \int_s^t F(X_r) - F(X_s) dX_r \\ &= F(X_s) X_{st} + DF(X_s) \underbrace{\int_s^t X_{sr} dX_r}_{=: \mathbb{X}_{st}} + R_{st} \end{aligned}$$

with  $\|R\|_{3\alpha} \leq C \|D^2 F\|_\infty \|X\|_\alpha^2$ , the integral becomes continuous w.r.t.  $(X_t, \mathbb{X}_{st})$  in  $\|X\|_\alpha + \|\mathbb{X}\|_{2\alpha}$ , sufficient for  $\alpha > \frac{1}{3}$ , since then

$$\int_s^t F(X_r) dX_r := \lim_{|\mathcal{P}| \rightarrow 0} \sum_{[u,v] \in \mathcal{P}} F(X_u) X_{uv} + DF(X_u) \mathbb{X}_{uv} \quad \text{well-defined}$$

- ▶ admissible integrands:  $Y \in C^\alpha$  s.th.  $Y_{st} = Y'_s X_{st} + R_{st}$  with  $Y' \in C^\alpha$  and  $R \in C^{2\alpha}$

## Back to Stochastic Filtering

$$\begin{aligned} E^{t,x} [g(X_T) | \mathcal{W}_{t,T}] &= \frac{E_0^{t,x} \left[ g(X_T) \exp \left( \int_t^T \gamma(X_s) dW_s - \frac{1}{2} \int_t^T \|\gamma(X_s)\|^2 ds \right) \right]}{E_0^{t,x} \left[ \exp \left( \int_t^T \gamma(X_s) dW_s - \frac{1}{2} \int_t^T \|\gamma(X_s)\|^2 ds \right) \right]} \\ &= \frac{u(t, x)}{u_0(t, x)} \end{aligned}$$

where

$$u(t, x) = g(x) + \int_t^T Lu(r, x) dr + \int_t^T \gamma_k u(r, x) dW_r^k$$

(and  $u_0$  the same backward eq with terminal condition  $g(x) \equiv 1$ )  
with

$$Lu := \frac{1}{2} \text{Tr} \left[ \sigma \sigma^T(x) D^2 u \right] + \langle b(x), Du \rangle$$

formally differentiating  $W^k$  yields the (backward rough) PDE

$$\begin{cases} -\partial_t u(t, x) &= Lu(t, x) + \gamma_k(x) u(t, x) \dot{W}_t^k \\ u(T, x) &= g(x), t \in [0, T], x \in \mathbb{R}^d. \end{cases}$$

# Rough linear PDE

$$\begin{cases} \partial_t u(t, x) &= Lu(t, x) + \Gamma_k u(t, x) \dot{W}_t^k \\ u(0, x) &= g(x), t \in [0, T], x \in \mathbb{R}^d. \end{cases} \quad (1)$$

where  $\mathbf{W}$  is a (geometric)  $\alpha$ -Hölder rough path,  $\alpha \in (1/3, 1/2]$  and

$$\begin{aligned} Lu &:= \frac{1}{2} \text{Tr} \left[ \sigma \sigma^T(x) D^2 u \right] + \langle b(x), Du \rangle + c(x)u(x) \\ \Gamma_k u &:= \langle \beta_k(x), Du \rangle + \gamma_k(x) u \end{aligned}$$

**Ansatz** in [Friz, Diehl, S. 2017]: (rough path) Feynman-Kac representation

$$u^{\mathbf{W}}(t, x) = \mathbb{E}^x \left[ g(X_t) \exp \left( \int_0^t c(X_r) dr + \int_0^t \gamma(X_r) \dot{W}_r dr \right) \right] \quad (2)$$

with

$$dX_t = \sigma(X_t) dB_t + b(X_t) dt + \beta(X_t) \dot{W}_t dt$$

where  $B$  is a Brownian motion

## Rough linear PDE, ctd.

$$\begin{cases} \partial_t u(t, x) &= Lu(t, x) + \Gamma_k u(t, x) \dot{W}_t^k \\ u(0, x) &= g(x), t \in [0, T], x \in \mathbb{R}^d. \end{cases} \quad (1)$$

where  $\mathbf{W}$  is a (geometric)  $\alpha$ -Hölder rough path,  $\alpha \in (1/3, 1/2]$  and

$$\begin{aligned} Lu &:= \frac{1}{2} \text{Tr} \left[ \sigma \sigma^T(x) D^2 u \right] + \langle b(x), Du \rangle + c(x)u(x) \\ \Gamma_k u &:= \langle \beta_k(x), Du \rangle + \gamma_k(x) u \end{aligned}$$

**Ansatz** in [Friz, Diehl, S. 2017]: understand how

$$u^{\mathbf{W}}(t, x) = \mathbb{E}^x \left[ g(X_t) \exp \left( \int_0^t c(X_r) dr + \int_0^t \gamma(X_r) \dot{W}_r dr \right) \right] \quad (2)$$

depends on  $\mathbf{W} = (W, \mathbb{W})$  in the **rough path metric**,

reduces to understand stability of  $\mathbf{W} \mapsto X$ , where

$$dX_t = \sigma(X_t) dB_t + b(X_t) dt + \beta(X_t) \dot{W}_t dt$$

# Rough linear PDE: weak solutions

## Theorem

[Friz, Nilssen, S. 2018]

Assume  $\sigma_{i,k}, \beta_j \in C_b^3(\mathbb{R}^d)$ ,  $b_j, \gamma_j, c \in C_b^1(\mathbb{R}^d)$ . Given  $g \in L^p(\mathbb{R}^d)$ , the Feynman-Kac integral (2) yields an analytically weak solution  $u$  of (1) satisfying

$$\sup_{t \in [0, T]} \|u(t)\|_{L^p} \leq C \|g\|_{L^p}$$

where  $C$  depends on  $\mathbf{W}$  only through the rough path metric.

**Rem** In the simple case  $b = c = \gamma = 0$

- ▶  $C = \mathbb{E} \left[ \sup_{x \in \mathbb{R}} |\det D\Phi_{0,t}^{-1}(x)| \right]$
- ▶  $\Phi$  being the flow generated by  $dX_t = \sigma(X_t) dB_t + \beta(X_t) d\mathbf{W}_t$



# Rough linear PDE: regular solutions

## Theorem

[Friz, Nilssen, S. 2018]

Assume  $\sigma_{i,k}, \beta_j, \gamma_j \in C_b^6(\mathbb{R}^d)$ ,  $b_j, c \in C_b^4(\mathbb{R}^d)$ . Given  $g \in W^{3,p}(\mathbb{R}^d)$ , the Feynman-Kac integral (2) yields a weak solution  $u \in W^{3,p}$  of (1) satisfying

$$\sup_{t \in [0, T]} \|u(t)\|_{W^{3,p}} \leq C \|g\|_{W^{3,p}}$$

where  $C$  depends on  $\mathbf{W}$  only through the rough path metric.

**Rem** In this case  $x \mapsto \Phi_{0,t}(x)$  is  $C_b^3(\mathbb{R})$ .

# Rough semilinear PDE

Extensions to the semilinear case

$$\begin{cases} \partial_t u(t, x) &= Lu(t, x) + F(u)(t, x) + \Gamma_k u(t, x) \dot{W}_t^k \\ u(0, x) &= g(x), t \in [0, T], x \in \mathbb{R}^d. \end{cases} \quad (3)$$

using Duhamels principle

$$u(t) = P_{0,t}^W g + \int_0^t P_{s,t}^W F(u(s)) ds \quad (\text{e.g.}) \text{ in } L^2$$

where  $P_{s,t}^W$  is the  $L^2$ -propagator of the linear Pb

requires a priori energy estimates of the type

- ▶  $\|u(t)\|_{L^2} \leq C \|g\|_{L^2}$
- ▶  $\int_0^t \|u(r)\|_{W^{1,2}}^2 dr \leq C \|g\|_{L^2}^2$

# Linear rough PDE: energy estimates

**Ansatz** Consider the rough heat equation

$$\begin{cases} \partial_t u(t, x) &= \Delta u(t, x) + \Gamma_k u(t, x) \dot{W}_t^k \\ u(0, x) &= g(x), t \in [0, T], x \in \mathbb{R}^d. \end{cases}$$

for smooth  $\mathbf{W}$

usual integration by parts leads to

$$\|u(t)\|_{L^2}^2 + 2 \int_0^t \|\nabla u(r)\|_{L^2}^2 dr = \|u(0)\|_{L^2}^2 + 2 \int_0^t (u(r))^2, \Gamma_k^* 1) dW_r^k$$

major problem: the elementary control of the 'rough integral'

$$\int_0^t (u(r))^2, \Gamma_k^* 1) dW_r^k \leq \|\Gamma_k^* 1\|_{L^\infty} \int_0^t \|u(r)\|_{L^2}^2 d|W_r^k|$$

is not continuous in the rough path metric

# Rough linear PDE: energy estimates, ctd.

Introducing the backward rough pde

$$\partial_t \varphi^{\mathbf{W}}(t, x) + \Delta \varphi^{\mathbf{W}}(t, x) + \Gamma_k^* \nabla \varphi^{\mathbf{W}}(t, x) \dot{W}_t^k + \Gamma_k^* 1(x) \varphi^{\mathbf{W}}(t, x) \dot{W}_t^k = 0$$

with nonnegative terminal condition, e.g.  $\varphi^{\mathbf{W}}(T, x) \equiv 1$  leads to

$$\int u(t)^2 \varphi^{\mathbf{W}}(t) dx + 2 \int_0^t \int |\nabla u(r)|^2 \varphi^{\mathbf{W}}(r) dx dr = \int u(0)^2 \varphi^{\mathbf{W}}(0) dx \quad t \in [0, T]$$

yields a priori energy estimate

as long as  $\mathbf{W} \mapsto \varphi^{\mathbf{W}}(t)$  is cont. in the rough path metric  
(e.g. via Feynman-Kac representation)

# Rough semilinear PDEs - main result

## Theorem

[Friz, Nilsson, S. 2018]

Assume  $\sigma_{i,k}, \beta_j, \gamma_j \in C_b^6(\mathbb{R}^d)$ ,  $b_j, c \in C_b^4(\mathbb{R}^d)$  as well as

- ▶  $\lambda|\xi|^2 \leq |\sigma^T \xi|^2$  for some  $\lambda > 0$ ,
- ▶  $\|F(u) - F(v)\|_{L^2} \leq C\|u - v\|_{W^{1,2}(\mathbb{R}^d)}$ .

Given  $g \in L^2(\mathbb{R}^d)$ , there exists a unique weak solution  $u \in C([0, T]; L^2) \cap L^2([0, T]; W^{1,2})$  of (3).

# Burgers equation

Navier Stokes equations

$$\begin{aligned}\partial_t u + u \cdot \nabla u + \nabla p &= \nu \Delta u, & (t, x) \in [0, T] \times \mathbb{R}^d \\ \operatorname{div} u &= 0, & u(0) = u_0 \in L^2(\mathbb{R}^d; \mathbb{R}^d).\end{aligned}$$

have no meaning when  $d = 1$

[Burgers, 1974] introduced the 1d toy-equation

$$\partial_t u + u \partial_x u = \nu \partial_x^2 u, \quad (t, x) \in [0, T] \times \mathbb{R}$$

neglecting the pressure in the Navier-Stokes equations.

# Stochastic Burgers equation

Adding randomness to the equation to incorporate

- ▶ fluctuations in classical continuum limits
- ▶ incorporate highly oscillating terms on small time-scales as statistical model for turbulence

first approach: additive, space-time white noise,  $\xi$

$$\partial_t u = \nu \partial_x^2 u - u \partial_x u + \xi,$$

- ▶ [Bertini, Cancrini, Jona-Lasinio 1994]
- ▶ [DaPrato, Debussche, Temam 1994] (1st step towards regularity structures)

# Stochastic Burgers equation with transport noise

Lagrangian viewpoint

$\phi_t(x) :=$  position of a fluid particle at time  $t$  starting in  $x$

assume the following decomposition

$$\dot{\phi}_t(x) = u(t, \phi_t(x)) - \beta_j(\phi_t(x)) \dot{Z}_t^j, \quad \phi_0(x) = x,$$

with  $(Z^j)_j$  modelling a highly oscillating part

associated velocity field  $u$  is generated by

$$\partial_t u = \nu \partial_x^2 u - u \partial_x u + \beta_j \partial_x u \dot{Z}^j \quad (4)$$

**Ansatz** [Hocquet, Nilssen, S. 2018] prove stability of

$(Z, \mathbb{Z}) \mapsto u$  w.r.t. rough path metric



# Variational approach

consider the weak formulation

$$\partial_t \int u_t \phi \, dx = - \int \nabla u_t \cdot \nabla \phi \, dx - \int u_t \operatorname{div}(\beta_j \phi) \, dx \dot{Z}_t^j$$

for all  $\phi \in W^{1,2}(\mathbb{R})$  w.r.t. the Gelfand triple

$$W^{1,2}(\mathbb{R}) \subset L^2(\mathbb{R}) \subset W^{-1,2}(\mathbb{R})$$

testing  $u$  against itself yields the energy estimate

$$\begin{aligned} \|u_t\|_{L^2}^2 + 2 \int_0^t \|\nabla u_r\|_{L^2}^2 \, dr &= \|u_0\|_{L^2}^2 - \int_0^t (u_r^2, \operatorname{div}(\beta_j)) \, dZ_r^j \\ &\leq \|u_0\|_{L^2}^2 + \|\operatorname{div}(\beta_j)\|_{L^\infty} \int_0^t \|u_r\|_{L^2}^2 \, d|Z^j|_r \end{aligned}$$

because of  $\int u_t \partial_x u_t \, dx = -\frac{1}{3} \int \partial_x (u_t^3) \, dx = 0$

## Rem

- ▶ again, no easy a priori (energy) estimates
- ▶ but also no easy perturbative approach

# Davies' method for solving rough ode's

**main idea** consider the rough ode

$$dy_t = g(y_t)dX_t,$$

a Taylor expansion yields

$$y_{st} := y_t - y_s = g(y_s)X_{st} + \nabla g(y_s)g(y_s)\mathbb{X}_{st} + y_{st}^{\natural},$$

**Definition**  $y$  is called a solution if  $y^{\natural}$  is more regular in time in the sense that

$$|y_{st}^{\natural}| \lesssim |t - s|^{1+}$$

## Davies' method for solving rough pde's

consider pure transport eqn (and a smooth path  $Z$ )

$$\partial_t f = (\beta_j \cdot \nabla) f \dot{Z}_t^j, \quad f|_{t=0} \in L^2(\mathbb{R}^d)$$

a similar Taylor expansion now yields

$$f_{st} = (\beta_j \cdot \nabla) f_s Z_{st}^j + (\beta_j \cdot \nabla)(\beta_k \cdot \nabla) f_s Z_{st}^{k,j} + f_{st}^{\text{h}}$$

where

$$f_{st}^{\text{h}}(x) = \int_s^t \int_s^{r_1} \int_s^{r_2} (\beta_j \cdot \nabla)(\beta_k \cdot \nabla)(\beta_i \cdot \nabla) f_{r_3}(x) dZ_{r_3}^i dZ_{r_2}^k dZ_{r_1}^j$$

**Rem** requires 3 derivatives of  $f$

# Davies' method for solving rough pde's, ctd.

for the variational approach therefore need

$$W^{3,2}(\mathbb{R}^d) \subset L^2(\mathbb{R}^d) \subset W^{-3,2}(\mathbb{R}^d)$$

define the unbounded rough drivers

$$A_{st}^1 \phi := (\beta_j \cdot \nabla) \phi Z_{st}^j \quad A_{st}^2 \phi := (\beta_j \cdot \nabla)(\beta_k \cdot \nabla) \phi Z_{st}^{k,j}$$

satisfying

$$|A_{st}^1|_{\mathcal{L}(W^n, W^{n-1})} \lesssim |t-s|^\alpha \quad |A_{st}^2|_{\mathcal{L}(W^n, W^{n-2})} \lesssim |t-s|^{2\alpha}$$

and Chen's relation  $A_{st}^2 - A_{s\theta}^2 - A_{\theta t}^2 = A_{\theta t}^1 A_{s\theta}^1$

**Definition**  $f : [0, T] \rightarrow L^2(\mathbb{R}^d)$  is called a solution if

$$f_{st} = A_{st}^1 f_s + A_{st}^2 f_s + f_{st}^{\natural}$$

with

$$|f_{st}^{\natural}(\phi)| \lesssim |t-s|^{1+} |\phi|_{W^3}.$$

# A priori estimates

consider the equation perturbed by a drift

$$f_{st} = \mu_{st} + A_{st}^1 f_s + A_{st}^2 f_s + f_{st}^\natural$$

where  $\mu : [0, T] \rightarrow W^{-2,2}(\mathbb{R}^d)$  is Lipschitz, e.g.  $\mu = \Delta f$

## Proposition

[Deya, Gubinelli, Hofmanova, Tindel 2016]

Suppose we have a solution to the above equation. Then

$$|f_{st}^\natural(\phi)| \lesssim (\mu_{Lip} \|A\|_\alpha + |f|_{L_{T,x}^{\infty,2}} \|A\|_\alpha^2) |\phi|_{W^3} |t - s|^{1+}.$$

# Stability

to obtain stability of

$$v_{st} = \int_s^t \Delta v_r dr + A_{st}^1 v_s + A_{st}^2 v_s + v_{st}^{\natural}$$

in the rough path topology we now need 3 derivatives

using a tensorization argument and involved approximation by smooth solutions we obtain

$$\begin{aligned} v_{st}^2(\phi) = & -2 \int_s^t |\nabla v_r|^2(\phi) + \nabla v_r(v_r \nabla \phi) dr \\ & + v_s^2(A_{st}^{1,*} \phi) + v_s^2(A_{st}^{2,*} \phi) + v_{st}^{2,\natural}(\phi) \end{aligned}$$

with the a priori estimate

$$|v_{st}^{2,\natural}(\phi)| \lesssim_A (\|\nabla v\|_{L^2([s,t] \times \mathbb{R}^d)}^2 + \|v\|_{L^\infty([s,t]; L^2(D))}^2) \|\phi\|_{W^{3,\infty}} |t-s|^{1+}$$

## Stability, ctd.

to obtain stability of

$$v_{st} = \int_s^t \Delta v_r dr + A_{st}^1 v_s + A_{st}^2 v_s + v_{st}^{\natural}$$

in the rough path topology we now need 3 derivatives

letting  $\phi = 1$  we get

$$\begin{aligned} (\|v\|_{L^2}^2)_{st} + 2 \int_s^t \|\nabla v_r\|_{L^2}^2 dr &= v_s^2(A_{st}^{1,*} 1) + v_s^2(A_{st}^{2,*} 1) + v_{st}^{2,\natural}(1) \\ &\lesssim \|A\|_{\alpha} |t - s|^{\alpha} (\|\nabla v\|_{L^2([s,t] \times \mathbb{R}^d)}^2 + \|v\|_{L^\infty([s,t]; L^2(\mathbb{R}^d))}^2) \end{aligned}$$

(rough) Gronwall lemma  $\Rightarrow$  energy estimate

## Application to Rough Burgers

$$\partial_t u = \partial_x^2 u - u \partial_x u + \beta_j \partial_x u \dot{Z}_t^j$$

formally yields the following eqn for  $u^2$

$$\partial_t u^2 = 2u \partial_x^2 u - 2u^2 \partial_x u + \beta_j \partial_x u^2 \dot{Z}_t^j$$

the remainder in this case can be estimated by

$$|u_{st}^{2,h}(\phi)| \lesssim_A \left( \|\nabla u\|_{L^2([s,t];L^2)}^2 + \|u\|_{L^\infty([s,t];L^2)}^2 + \|u\|_{L^\infty([s,t];L^3)}^3 \right) \cdot \|\phi\|_{W^{3,\infty}} |t - s|^{1+}$$

**Rem** cubic in  $u$ , so rough Gronwall lemma not applicable



## Weighted measure space approach

similar to part 1 introduce the backward rough pde

$$\partial_t \varphi_t^Z = -\partial_x^2 \varphi_t^Z + \partial_x(\beta_j \varphi_t^Z) \dot{Z}_t^j, \quad \varphi_T^Z \equiv 1$$

leads to the weighted energy equality

$$\partial_t \int u_t^2 \varphi_t^Z dx + 2 \int (\partial_x u_t)^2 \varphi_t^Z dx = 2 \int u_t^2 \partial_x u_t \varphi_t^Z dx$$

if we can show that  $0 < m_* \leq \varphi_t^Z(x) \leq m^* < \infty$

$$\begin{aligned} \|u_T\|_{L^2}^2 + \int_0^T \|\partial_x u_t\|_{L^2}^2 dt \\ \lesssim_{m_*, m^*} \|u_0\|_{L^2}^2 + 2 \int_0^T (|u_t|^2, |\partial_x u_t|) dt \end{aligned}$$

still cubic in  $u$  on RHS

Use non-linear Gronwall:

### Theorem (Bihari-LaSalle)

Suppose  $q > 1$  and two positive functions  $x, k : [0, T] \rightarrow \mathbb{R}_+$  satisfy

$$x_t \leq x_0 + \int_0^t k_s x_s^q ds.$$

Then for all  $t$  such that  $(q - 1)x_0 \int_0^t k_s ds < 1$  we have the following estimate

$$x_t \leq \frac{x_0}{\left(1 - (q - 1)x_0 \int_0^t k_s ds\right)^{1/(q-1)}}.$$

# Backward equation

main advantage of the backward equation

$$\partial_t \varphi_t^Z = -\partial_x^2 \varphi_t^Z + \partial_x (\beta_j \varphi_t^Z) \dot{Z}_t^j, \quad m_T = 1$$

admits the Feynman-Kac representation

$$\varphi_t^Z(x) = E \left[ \exp \left( \int_t^T (\partial_x \beta_j)(\psi_{t,s}(x)) dZ_s^j \right) \right]$$

where

$$d\psi_{t,s}(x) = \sqrt{2} dB_s + \beta_j(\psi_{t,s}(x)) dZ_s^j, \quad \psi_{t,t}(x) = x$$

So in fact this method can be seen as a flow-transformation on the energy space

## Main result

This gives

$$\partial_t(u_t^2, m_t) + 2(|\partial_x u_t|^2, m_t) = 2(u_t^2, \partial_x(u_t m_t)).$$

Upper and lower bound on  $m$  can be checked directly from the Feynman-Kac representation.

## Theorem

[Hocquet, Nilssen, S. 2018]

There exists a maximal time  $T_0 \in (0, T]$  such that existence and uniqueness of finite-energy solutions  $u$  to

$$\partial_t u = \partial_x^2 u - u \partial_x u - \beta_j \partial_x u \dot{Z}^j$$

holds on  $(0, T_0) \times \mathbb{R}$

The value of  $T_0$  depends only upon the quantities  $\|u_0\|_{L^2}$ ,  $\|\beta_j\|_{C_b^6}$ ,  $\|Z\|_\alpha$ ,  $\|\mathbb{Z}\|_{2\alpha}$